

HOMOGENEOUS AND ISOTROPIC
STATISTICAL SOLUTIONS
OF THE NAVIER-STOKES EQUATIONS

Andrei V. Fursikov

Department of Mechanics and Mathematics
Moscow State University, 119991 Moscow,
Russia

Stochastics in Turbulence and Finance
Sonderborg, January 30, 2008

The aim of the talk is to expound results on Isotropic Statistical Solution from

S.Dostoglou, A.V.Fursikov, J.D.Kahl: *Homogeneous and Isotropic Statistical Solutions of the Navier-Stokes Equations.*- Math. Physics Electronic Journal,
<http://www.ma.utexas.edu/mpej/> volume 12,
paper No. 2, 2006

These results are founded on results on Homogeneous Statistical Solutions published in

Vishik, V.I. and A.V. Fursikov: *Mathematical problems of statistical hydromechanics.* Kluwer, 1988.

Initial definitions

Let H be a Hilbert space of vector fields $u(x) : u = (u_1, u_2, u_3), x = (x_1, x_2, x_3) \in \mathbb{R}^3$.

$T_h : H \rightarrow H, h \in \mathbb{R}^3$ is translation operator if:

$$T_h u(x) = u(x + h).$$

$O(3)$ denotes the group of all orthogonal 3×3 -matrices ω (with $\det \omega = \pm 1$)

$R_\omega : H \rightarrow H, \omega \in O(3)$ is called rotation operator if

$$(R_\omega u)(x) = \omega u(\omega^{-1}x)$$

In fact R_ω can be also reflection operator for some ω

Homogeneous and isotropic measures

$\mathcal{B}(H)$ is the σ -algebra of Borel sets of H .

Let H_1, H_2 be Hilbert spaces. For a measure $\mu(A), A \in \mathcal{B}(H_1)$, **push forward of μ under the map $\Psi : H_1 \rightarrow H_2$** is the measure

$$\Psi^* \mu(B) = \mu(\Psi^{-1}B) \quad \forall B \in \mathcal{B}(H_2).$$

where

$$\Psi^{-1}B := \{m \in H_1 : \Psi(m) \in B\} \in \mathcal{B}(H_1).$$

A measure μ defined on $\mathcal{B}(H)$ is called **homogeneous** if it is translation invariant:

$$T_h^* \mu(A) = \mu(A), \quad \forall A \in \mathcal{B}(H), h \in \mathbb{R}^3 \quad (1)$$

A homogeneous measure μ defined on $\mathcal{B}(H)$ is called **isotropic** if

$$R_\omega^* \mu(A) = \mu(A), \quad \forall A \in \mathcal{B}(H), \omega \in O(3) \quad (2)$$

Why isotropic measures are usefull?

Correlation tensor for general (nonisotropic) measure is defined as follows:

$$K_{i,j}(x_1, x_2) = \int u(x_1) \otimes u(x_2) \mu(du),$$

where $i, j = 1, \dots, 3$, $x_1, x_2 \in \mathbb{R}^3$.

If μ is isotropic measure, its correlation tensor has the form:

$$K_{i,j}(x_1, x_2) = k(|x_1 - x_2|) \delta_{i,j}$$

where $k(\lambda)$, $\lambda \in \mathbb{R}_+$ is a scalar function, and $\delta_{i,j}$ is Kroneker symbol.

In Kolmogorov turbulence theory isotropic measures are essentially used.

Navier-Stokes Equations

$$\partial_t u(t, x) - \Delta u + (u, \nabla u) + \nabla p(t, x) = 0, \quad \operatorname{div} u = 0,$$

$$u(t, x)|_{t=0} = u_0(x)$$

where $t \in (0, T)$, $x \in \mathbb{R}^3$.

Energy inequality

$$\|u(t, \cdot)\|^2 + \int_0^t \|\nabla u(\tau, \cdot)\|^2 d\tau \leq \|u_0\|^2, \quad t \in (0, T),$$

($\|\cdot\| = \|\cdot\|_{L_2(\mathbb{R}^3)}$) is the basic property to prove existence of generalized solution.

Energy inequality can be used to prove existence of statistical solution that is, roughly speaking, the measure P supported on the set of generalized solution and satisfying the following analog of energy estimate:

$$\int (\|u(t, \cdot)\|^2 + \int_0^t \|\nabla u(\tau, \cdot)\|^2 d\tau) P(du) \leq \int \|u_0\|^2 \mu(du_0)$$

where μ is a given measure on the set of initial conditions.

Properties of homogeneous measures

1) Unique homogeneous measure supported on $L_2(\mathbb{R}^3)$ is δ -measure, i.e. $\delta(A) = 1$ for $A \in \mathcal{B}(L_2(\mathbb{R}^3))$ if $0 \in A$ and $\delta(A) = 0$ if $0 \notin A$

2) Each homogeneous measure μ is supported on bounded functions defined on \mathbb{R}^3 . Moreover, if homogeneous measure $\mu \neq \delta$ (and does not contain component $\alpha\delta, \alpha \in (0, 1)$), then $\mu(L_2(\mathbb{R}^3)) = 0$.

That is why the Hilbert space H on which there is reason to consider homogeneous (and isotropic) measures is as follows:

$$\mathcal{H}^0(r) = \{u(x), x \in \mathbb{R}^3 : \operatorname{div} u = 0,$$

$$\|u\|_{\mathcal{H}^0(r)}^2 = \int (1 + |x|^2)^r |u(x)|^2 dx < \infty\}$$

where $r < -3/2$.

Since $\|u\|_{L_2(\mathbb{R}^3)} = \infty$ for typical $u \in \mathcal{H}^0(r)$, there is no hope to use energy estimate to prove existence of homogeneous statistical solution for Navier-Stokes equation.

Density of energy for homogeneous measures

For a homogeneous measure μ the pointwise averages

$$\int |u|^2(x) \mu(du), \quad \int |\nabla u|^2(x) \mu(du) \quad (3)$$

can be defined by the equalities

$$\int \int |u(x)|^2 \phi(x) dx \mu(du) = \int |u(x)|^2 \mu(du) \int \phi(x) dx$$

$$\int \int |\nabla u(x)|^2 \phi(x) dx \mu(du)$$

$$= \int |\nabla u(x)|^2 \mu(du) \int \phi(x) dx \quad \forall \phi \in L_1(\mathbb{R}^3)$$

The first expression in (3) is **the energy density** and the second one is **the density of the energy dissipation**. In terms of these qualities we will get analog of energy estimate for statistical solution.

Some definitions

The set of all **generalized solutions** of the Navier-Stokes system:

$$\mathcal{G}_{NS} = \left\{ u \in L^2(0, T; \mathcal{H}^0(r)) : L(u, \phi) \right. \\ \left. \equiv \int_0^T \left(\langle u, \frac{\partial \phi}{\partial t} \rangle_2 + \langle u, \Delta \phi \rangle_2 + \sum_{j=1}^3 \langle u_j u, \frac{\partial \phi}{\partial x_j} \rangle_2 \right) dt = 0, \right.$$

for all $\phi \in C_0^\infty((0, T) \times \mathbb{R}^3) \cap C((0, T); \mathcal{H}^0(r))$,
 where $\langle u, v \rangle_2 = \int_{\mathbb{R}^3} u(x) \cdot v(x) dx$.

A measure $P(A)$, $A \in \mathcal{B}(L^2(0, T; \mathcal{H}^0(r)))$ is called **homogeneous in x** if $T_h^* P = P$

$\forall h \in \mathbb{R}^3$

The space $\mathcal{H}^1(r)$ is defined as follows:

$$\mathcal{H}^1(r) = \{ u(x) \in \mathcal{H}^0(r) : \|u\|_{\mathcal{H}^1(r)} \equiv \|\nabla u\|_{\mathcal{H}^0(r)}^0 \leq \infty \}$$

Definition of homogeneous statistical solution

Given homogeneous probability measure μ on $\mathcal{B}(\mathcal{H}^0(r))$ possessing finite energy density, a **homogeneous statistical solution of the Navier-Stokes equations with initial condition** μ is a probability measure P on $\mathcal{B}(L^2(0, T; \mathcal{H}^0(r)))$ such that:

1. P is homogeneous in x .
2. $P(\widehat{W}) = 1$, where $\widehat{W} = L^2(0, T; \mathcal{H}^1(r)) \cap BV^{-s} \cap \mathcal{G}_{NS}$, $s > \frac{11}{2}$.
3. $P(\gamma_0^{-1}A) = \mu(A)$, $\forall A \in \mathcal{B}(\mathcal{H}^0(r))$
 where $\gamma_0^{-1}A = \{u \in \widehat{W} : \gamma_0 u \in A\}$.
4. $\int \left(|u(t, x)|^2 + \int_0^t |\nabla u|^2(\tau, x) d\tau \right) P(du) \leq C \int |u(x)|^2 \mu(du)$,
 for each $t \in [0, T]$.

Definition of isotropic statistical solution

A homogeneous measure $P(A)$, $A \in \mathcal{B}(L^2(0, T; \mathcal{H}^0(r)))$ is called **isotropic in x** if $R_\omega^* P = P, \forall \omega \in O(3)$

Let μ be an isotropic measure on $\mathcal{B}(\mathcal{H}^0(r))$ possessing finite energy density.

An **isotropic statistical solution of the Navier-Stokes equations with initial condition μ** is a homogeneous statistical solution P that is isotropic in x

The main results

Theorem 1 Given μ homogeneous measure on $\mathcal{H}^0(r)$ with finite energy density, there exists homogeneous statistical solution of the Navier-Stokes equations P with initial condition μ .

Theorem 2 Given μ isotropic measure on $\mathcal{H}^0(r)$ with finite energy density, there exists isotropic statistical solution of the Navier-Stokes equations \hat{P} with initial condition μ .

Draft of Theorem's 1 proof

Step 1. Introduce the finite-dimensional space of trigonometric polynomials of degree l and period $2l$:

$$\mathcal{M}_l = \left\{ \sum_{k \in \frac{\pi}{l}\mathbb{Z}^3, |k| \leq l} a_k e^{ik \cdot x} : a_k \cdot k = 0, a_k = \bar{a}_{-k} \forall k \right\},$$

For all $l \in \mathbb{N}$ $\mathcal{M}_l \subset \mathcal{H}^0(r)$.

Lemma 1. For a given initial homogeneous measure μ on $\mathcal{H}^0(r)$ with finite energy density there exists a sequence of homogeneous measures μ_l as $l \rightarrow \infty$ defined on $\mathcal{B}(\mathcal{H}^0(r))$ and supported on \mathcal{M}_l such that

$$\int \exp(i(u, \varphi)) \mu_l(du) \rightarrow \int \exp(i(u, \varphi)) \mu(du)$$

as $l \rightarrow \infty$, and

$$\int |u|^2(x) \mu_l(du) \leq \int |u|^2(x) \mu(du)$$

Step 2. Galerkin's approximations

$$\partial_t u(t, x) - \Delta u + \pi_l(u, \nabla)u = 0, \quad \operatorname{div} u = 0,$$

$$u(t, x) \in C(0, T; \mathcal{M}_l), \quad u|_{t=0} = u_0 \in \mathcal{M}_l$$

where π_l is the projection on \mathcal{M}_l . Let $S_l : \mathcal{M}_l \rightarrow C(0, T; \mathcal{M}_l)$ be operator mapping initial conditions to solutions of this problem (i.e. Galerkin's resolving operator).

Galerkin's approximations of homogeneous statistical solution for Navier-Stokes equations are the measures P_l that are defined with help of approximation μ_l from Lemma 1 of initial measure μ by the formula:

$$P_l(A) = \mu_l(S_l^{-1}A) \quad \forall A \in \mathcal{B}(C(0, T; \mathcal{M}_l))$$

Step 3. Passage to limit

$$P_l \rightarrow P \quad \text{weakly as } l \rightarrow \infty$$

on $L_2(0, T; \mathcal{H}^0(r))$ where P is homogeneous statistical solution of Navier-Stokes equations.

Isotropic statistical solution

Examples of isotropic measures: for homogeneous measure μ define

$$\hat{\mu}(A) = \int_{O(3)} R_{\omega}^* \mu(A) d\omega \quad \forall A \in \mathcal{B}(\mathcal{H}^0(r))$$

where $d\omega = H$ is the standard Haar measure on $O(3)$ normalized. Then $\hat{\mu}$ is isotropic measure.

Draft of Theorem's 2 proof

Let P be a homogeneous statistical solution constructed in Theorem 1. Then

$$\hat{P}(A) = \int_{O(3)} R_{\omega}^* P(A) d\omega \quad \forall A \in \mathcal{B}(\mathcal{H}^0(r))$$

is isotropic statistical solution of Navier-Stokes equations.

Approximations of isotropic statistical solution

Spaces for isotropic approximations:

$$\widehat{\mathcal{M}}_l = \bigcup_{\omega \in O(3)} R_\omega M(l)$$

where, recall,

$$\mathcal{M}_l = \left\{ \sum_{k \in \frac{\pi}{l}\mathbb{Z}^3, |k| \leq l} a_k e^{ik \cdot x} : a_k \cdot k = 0, a_k = \bar{a}_{-k} \forall k \right\},$$

Let $S_l : \mathcal{M}_l \rightarrow C(0, T; \mathcal{M}_l)$ be Galerkin's resolving operator. Given v in $\widehat{\mathcal{M}}_l$, there are ω in $O(3)$ and u in \mathcal{M}_l such that $v = R_\omega u$. Extend S_l from \mathcal{M}_l to $\widehat{\mathcal{M}}_l$ as

$$\widehat{S}_l v = R_\omega S_l u.$$

Lemma 2 This extension is well-defined.

Let μ be initial measure, μ_l be its periodic approximation,

$$\widehat{\mu}_l(A) = \int_{O(3)} R_\omega^* \mu_l(A)$$

be isotropic averaging of μ_l . We set

$$\widehat{P}_l(A) = \widehat{\mu}_l(\widehat{S}_l^{-1} A), \quad A \in \mathcal{B}(L^2(0, T, \mathcal{H}^0(r)))$$

Lemma 3 The measure \widehat{P}_l is homogeneous and isotropic. Moreover

$$\widehat{P}_l(A) = (P_l \times H)(a^{-1} A)$$

where $a(\omega, u) = R_\omega u$ and H is the Haar measure on $O(3)$, normalized.

Theorem 3 $\widehat{P}_l \rightarrow \widehat{P}$ weakly on $L_2(0, T; \mathcal{H}^0(r))$ as $l \rightarrow \infty$, where \widehat{P} is the isotropic statistical solution, constructed in Theorem 2.